

DETAILED STUDY OF NUCLEAR DENSITY AND CHARGE DISTRIBUTION EFFECTS IN CUMULATIVE PARTICLE PRODUCTION

V.K.Bondarev*, A.G.Litvinenko, P.I.Zarubin

Detailed studies of A-dependence of a cumulative particle production carried out on beams of protons and light nuclei have shown similarities in a behaviour of the cross sections per nucleon on different beams in a wide range of fragmenting nuclei. The cross sections of a cumulative particle production are closely correlated with a nuclear density in a region of light nuclei. There is a correlation of cross sections with a charge density distribution in a region of medium weight nuclei. The cross sections of positive and negative particles are correlated with a relative content of protons and neutrons in nuclei. A large neutron excess in heavy nuclei doesn't produce a difference of the cross sections for positive and negative pions (both are approximately equal).

The investigation has been performed at the Laboratory of High Energies, JINR.

Детальное изучение влияния ядерной плотности и распределения заряда на кумулятивное рождение частиц

В.К.Бондарев, А.Г.Литвиненко, П.И.Зарубин

Детальное изучение A-зависимости рождения кумулятивных частиц на пучках протонов и легких ядер показывает сходное поведение сечения на нуклон для разных пучков в широкой области фрагментирующих ядер. Сечение рождения кумулятивных частиц коррелирует с распределением ядерной плотности для легких ядер. Корреляция с плотностью заряда наблюдается в области ядер со средним весом. Сечение рождения положительных и отрицательных частиц коррелирует с относительным содержанием протонов и нейтронов в ядрах. Большой нейтронный избыток для тяжелых ядер не сказывается на сечении рождения положительных и отрицательных пионов (оба сечения примерно равны).

Работа выполнена в Лаборатории высоких энергий ОИЯИ.

An investigation of a cumulative particle production made it possible to establish general features of a target nucleus fragmentation in hadron-

*St.Petersburg University, Russia

nucleus and nucleus-nucleus collisions [1,2]. The most important one is a universality of an exponential slope parameter of scale variable spectra of cumulative particles in a wide range of collision energy and fragmenting nuclei. On the basis of these results a quark-parton structure function of nucleus was introduced [3], which was supposed to reflect a momentum quark distribution in nuclei. A behaviour of the structure function was studied in a range of a relativistic invariant scale variable $0.35 < X < 3.5$ [4]. A relativistic invariant definition of X can be found in [4].

Our previous results on A -dependence with proton beams [1,2,3] were confirmed in the paper [5].

A dominant feature in a behaviour of an A -dependence of cross sections of cumulative particles (π , K , p , bound systems like deuterons, tritons) is a rapid growth of cross sections in the region of the lightest nuclei which becomes more flat with a growth of a mass number of fragmenting nuclei. For pions this change in a growth corresponds to $A \approx 30$; and for protons and baryon fragments, to $A \approx 100$.

A fragmentation of the following nuclei was studied in [1,2,3,4,5]: D, ^4He , ^6Li , ^7Li , Be, C, Al, Si, ^{54}Fe , ^{56}Fe , ^{58}Fe , ^{58}Ni , ^{61}Ni , ^{64}Ni , ^{64}Zn , ^{114}Sn , ^{124}Sn , Pb. Separated isotope targets give an opportunity to study a cumulative production at a fixed charge of nuclei with a variation of a neutron content. On the other side, it makes possible to vary a charge of nucleus at a fixed mass number.

The dependence of a ratio of structure functions on a density of nuclei (^6Li , ^{12}C , ^{40}Ca) at X less than 0.06 was precisely measured by the NMC collaboration [6]. Our data were taken in a range of the scale variable $0.7 \leq X \leq 1.4$. To analyze density effects we used the same expression for a density of nuclei as in the paper [6] of the NMC collaboration:

$$\rho(A) = A / \left[\frac{4}{3}\pi \left(\frac{5}{3} \langle r^2 \rangle \right)^{3/2} \right],$$

where $\langle r^2 \rangle$ is the mean square nuclear charge radius [7].

Inclusive invariant cross sections of secondary particles (π^+ , π^- , p , d) normalized per nucleon of a fragmenting nucleus were measured at a production angle of 120° with a momentum of 0.5 GeV/c for collisions of $4.5A$ GeV/c deuterons with nuclei [5]. An expression for a cross section is shown below.

$$\frac{1}{A} \cdot E \frac{d\sigma}{dp} \equiv \frac{1}{A} \cdot \frac{E}{p^2} \frac{d^2\sigma}{dpd\Omega} \text{ (mb} \cdot \text{GeV}^{-2} \cdot \text{c}^3 \cdot \text{sr}^{-1} \cdot \text{nucleon}^{-1}\text{)}.$$

Figure 1 shows A -dependence of the cross section per nucleon in a region of the lightest nuclei. There is a definite correlation of the cross

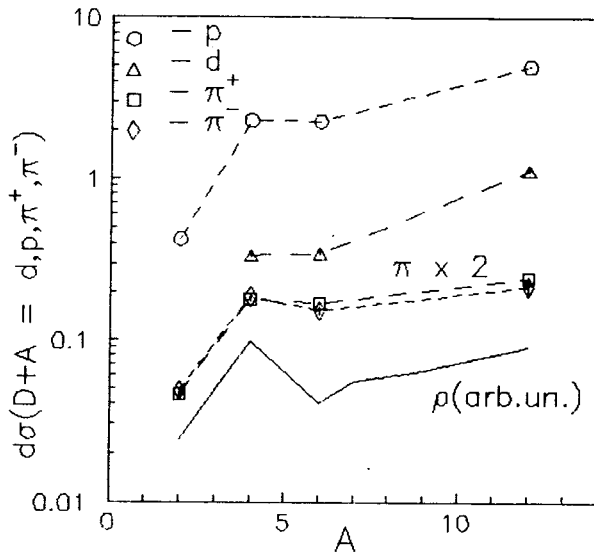


Fig.1. A-dependence of cross sections of a cumulative particle production in $D + A \rightarrow c + \dots$ ($\theta = 120^\circ$)

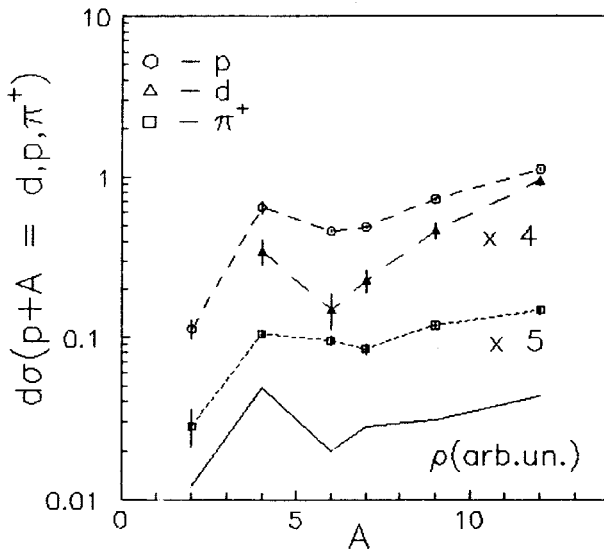


Fig.2. A-dependence of cross sections of a cumulative particle production in $p + A \rightarrow c + \dots$ ($\theta = 180^\circ$)

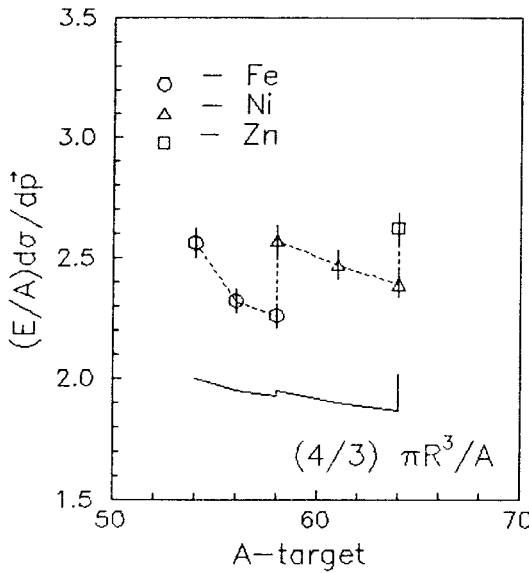


Fig.3. A-dependence of cross sections of cumulative particle production in $p + A \rightarrow p + \dots$ ($p_0 = 8.9 \text{ GeV/c}$, $\theta = 180^\circ$)

sections with a density of a fragmenting nucleus (lower curve). A ratio of helium to deuteron nuclear densities is equal to 4.2 and corresponds to a growth of a relative yield for protons by a factor of 5.38 ± 0.22 ; for π^+ , 3.90 ± 0.13 ; for π^- , 3.86 ± 0.13 . A density of ${}^6\text{Li}$ is 2.5 times lower than ${}^4\text{He}$, but cross sections per nucleon are approximately equal. A density of a carbon nucleus is 2.25 times higher than for ${}^6\text{Li}$ and corresponds to the ratio of cross sections 2.22 ± 0.06 for protons, 3.26 ± 0.09 for deuterons, 1.43 ± 0.04 for π^+ and 1.42 ± 0.17 for π^- .

Figure 2 shows a similar data for collisions of 8.9 GeV/c protons with nuclei [4,8]. Secondary particles were observed at an angle of 180° with momenta of 0.5 GeV/c. A correlation between a density (lower curve) and a cross section is more clear, especially for a transition from a helium nucleus to a lithium one.

As a mass number A increases, a density tends to a constant value, and an observation of such a correlation becomes less probable. Nevertheless, irregularities related with a fraction of a neutron and proton content were seen in this region [2].

It was experimentally shown that proton production cross sections are independent of a neutron excess at a fixed charge of a nucleus. This effect was called an isotopic effect. In paper [9], this conclusion was confirmed and an isotonic effect was established, i.e., an independence of a neutron yield on a number of protons in a nucleus. A detailed study of a cumulative particle production [5] on nuclei with a well defined number of protons and neutrons made it possible to obtain a qualitatively new information on these phenomena.

Proton production cross sections on isotopes of Fe, Ni, and Zn [8] are shown in fig.3. A primary beam is 8.9 GeV/c protons, a secondary proton

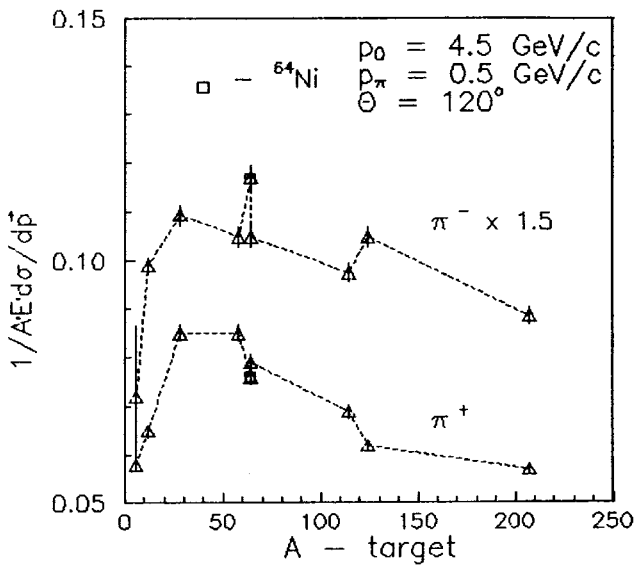


Fig.4. A-dependence of cross sections of cumulative particle production in $p + A \rightarrow p + \dots$

momentum is $0.5 \text{ GeV}/c$; a production angle, 180° . The experimental points are connected with a dashed line, and the lower curve fits the formula

$$V = \frac{4}{3} \pi R^3 / A,$$

which we assign as a charge density of a nucleus. $R = \langle r^2 \rangle^{1/2}$ is a mean radius of a nucleus. There is a correlation between the cross sections and this parameter, especially for Ni and Zn isotopes. The cross section shows a sharp growth at a small growth of a charge from ^{58}Fe to ^{58}Ni , from ^{64}Ni to

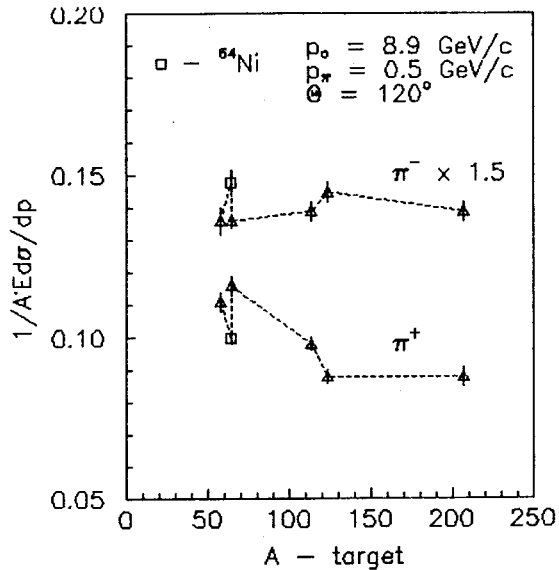


Fig.5. A-dependence of cross sections of a cumulative particle production in $p + A \rightarrow p + \dots$

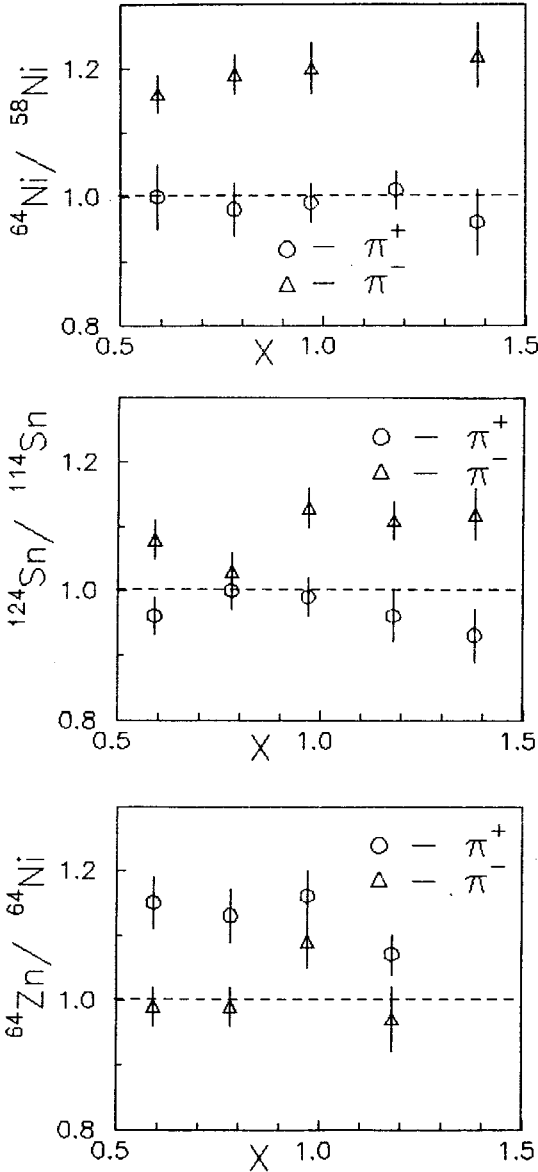


Fig.6. Ratio of cross sections for positive and negative pions

^{64}Zn ($Z = 26, 28, 30$). The results on π^+ and π^- production in collisions of 4.5 and 8.9 GeV/c protons with nuclei are shown in figs.4 and 5. Pions were observed with 0.5 GeV/c momentum at an angle of 120° . The isotopic effect can be seen on both figures.

We studied the isotopic effect for pions and protons in a range of secondary particles momenta from 0.3 up to 0.7 GeV/c. Figure 6 (pions) and fig.7 (protons) show ratio of cross sections vs scale variable X , corresponding to momenta of this range. It is supposed a ratio of cross sections is proportional to a ratio of quark-parton structure functions of nuclei. Let us describe it in detail.

Cross sections in a $^{64}\text{Ni}/^{58}\text{Ni}$ combination (same Z , different N) are equal for positive pions and follow to a neutron excess for negative ones. For a

$^{64}\text{Zn}/^{64}\text{Ni}$ combination we

observe an opposite situation. A behaviour of a $^{124}\text{Sn}/^{114}\text{Sn}$ combination is similar to the first one.

For protons (fig.7) a $^{64}\text{Ni}/^{58}\text{Ni}$ combination shows an equality of cross sections or an independence of cross sections on a neutron excess. In case of a $^{64}\text{Zn}/^{64}\text{Ni}$ combination a dominating role of a charge of a nucleus is seen

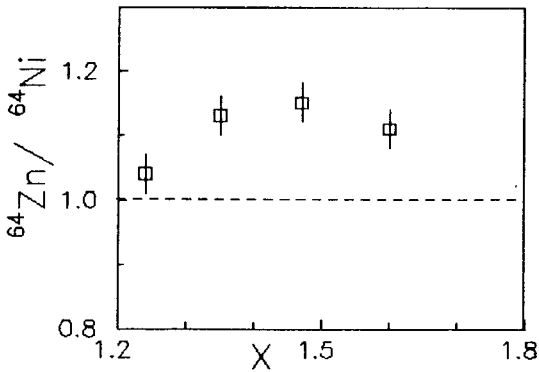
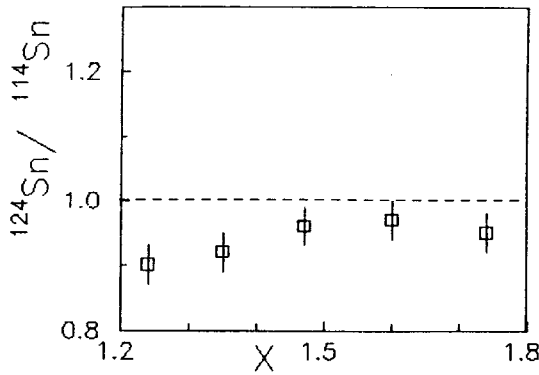
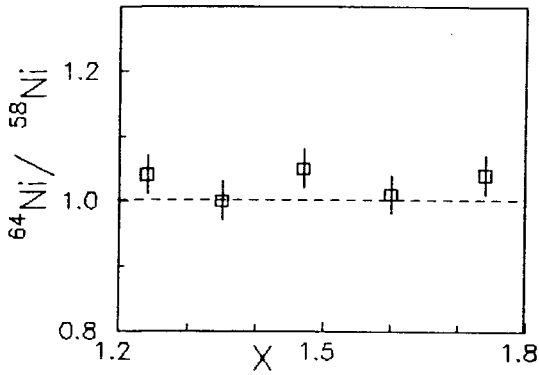


Fig.7. Ratio of cross sections for protons

for proton production. A $^{124}\text{Sn}/^{114}\text{Sn}$ ratio is different from a constant and less than unity.

Let us make a summary of our conclusions:

- the cross sections of cumulative particle production are closely correlated with a nuclear density in a region of light nuclei;

- there is a correlation of cross sections with a charge density distribution in a region of medium weight nuclei;

- the cross sections of positive and negative particles are correlated with a relative content of protons and neutrons in nuclei;

- a large neutron excess in heavy nuclei doesn't produce a difference of the cross sections for positive and negative pions (both are approximately equal);

- these conclusions are considered as additional test of models of the cumulative effect and, in general, models describing collisions of relativistic nuclei.

References

1. Baldin A.M. — Sov. J. Particles and Nuclei, 1977, 8(3), p.175 (publ.Amer. Inst. Phys.).

2. Stavinsky V.S. — Sov. J. Particles and Nuclei, 1979, 10(5), p.949.
3. Baldin A.M. — Proc. Intern. Conf. on Extreme States in Nuclear Systems. Dresden, 1980, 2, p.1.
4. Baldin A.M. et al. — JINR Commun. E1-82-472, Dubna, 1982.
5. Averichev G.S. et al. — In: Relativistic Nuclear Physics & Quantum Chromodynamics Proc. of X Intern. Seminar on High Energy Physics Problems (24—29 Sept. 1990, Dubna) pp.90—96 World Scientific (1991).
6. CERN — NMC P.Amaudruz et al. — NMC-collaboration, CERN-PPE/91-147, 9th Sept. 1991.
7. H. de Vries et al. — Atomic Data and Nuclear Data Tables, 1987, 36, p.495.
8. Baldin A.M. et al. — JINR Commun. R1-11302, Dubna, 1978.
9. Gavrilov V.B. et al. — Preprint ITEP-121 Moscow, 1985.

Received on December 28, 1992.